

$$1. A \rightarrow B + C$$

$$\quad \quad \quad \|\vec{p}_B\| \quad \|\vec{p}_C\|$$

Express E_B, E_C, p_B, p_C all in terms of m_A, m_B, m_C .

In the rest frame of A: $P_A^\mu = \begin{pmatrix} m_A c \\ \vec{0} \end{pmatrix}$ and since $P_A^\mu = P_B^\mu + P_C^\mu$

$$\text{we consider } P_A^\mu - P_B^\mu = P_C^\mu$$

It helps to recall that for a single particle

$$P_\mu P^\mu = -m^2 c^2$$

To see this recall that $P_\mu P^\mu$ is invariant

so we can just evaluate it in the particle's rest frame where $P^\mu = \begin{pmatrix} mc \\ \vec{0} \end{pmatrix} \Rightarrow P_\mu = (-mc, \vec{0})$

$$\text{then } (P_{A\mu} - P_{B\mu})(P_A^\mu - P_B^\mu) = P_{C\mu} P_C^\mu$$

$$P_{A\mu} P_A^\mu - 2P_{A\mu} P_B^\mu + P_{B\mu} P_B^\mu = P_{C\mu} P_C^\mu$$

$$-m_A^2 c^2 - 2m_A E_B - m_B^2 c^2 = -m_C^2 c^2$$

$$E_B = \frac{1}{2m_A} [-m_C^2 c^2 + m_A^2 c^2 + m_B^2 c^2]$$

and by inspection

$$E_C = \frac{1}{2m_A} [-m_B^2 c^2 + m_A^2 c^2 + m_C^2 c^2]$$

Then using that: $\frac{E^2}{c^2} - p^2 = m^2 c^2 \Rightarrow p^2 = \frac{E^2}{c^2} - m^2 c^2$

$$\text{we find: } p_B = \sqrt{\frac{E_B^2}{c^2} - m_B^2 c^2} = \sqrt{\frac{1}{4m_A^2 c^2} [-m_C^2 c^2 + m_A^2 c^2 + m_B^2 c^2]^2 - m_B^2 c^2}$$

$$\text{and: } p_C = \sqrt{\frac{E_C^2}{c^2} - m_C^2 c^2} = \sqrt{\frac{1}{4m_A^2 c^2} [-m_B^2 c^2 + m_A^2 c^2 + m_C^2 c^2]^2 - m_C^2 c^2}$$

It turns out that if you multiply these out, you will find that $p_B = p_C$. But this is as expected since $\vec{p}_{\text{tot}} = \vec{0}$ for the decay of a particle at rest.

2. Recall that $F_{\mu\nu} = \begin{pmatrix} 0 & -E_x & -E_y & -E_z \\ E_x & 0 & B_z & -B_y \\ E_y & -B_z & 0 & B_x \\ E_z & B_y & -B_x & 0 \end{pmatrix}$

so $F_{23} = B_x$ and $F_{32} = 0$, but we want to get these from $F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$ where $A_\mu = (\phi, \vec{A})$ and $\vec{E} = -\vec{\nabla}\phi - \dot{\vec{A}}$ and $\vec{B} = \vec{\nabla} \times \vec{A}$

Thus: $F_{23} = \partial_2 A_3 - \partial_3 A_2 = \frac{\partial A_z}{\partial y} - \frac{\partial A_y}{\partial z}$, but this is the x-component of $\vec{\nabla} \times \vec{A}$, hence B_x .
 $F_{12} = \partial_2 A_1 - \partial_1 A_2 = 0$ obviously

3. First of all $g_{R_{xy}} = \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & -\frac{1}{2} \end{pmatrix} = \frac{1}{2} \sigma_x$ where $\sigma_1 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \Rightarrow \sigma_2 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \Rightarrow \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \sigma_1$, etc.

Now note that $g_{R_{xy}}^2 = \frac{1}{4} \sigma_x^2 = \frac{1}{4} I$ and $g_{R_{xy}}^3 = \frac{1}{8} \sigma_x^3 = \frac{1}{8} \sigma_x$, in general $g_{R_{xy}}^n = \begin{cases} \frac{1}{2^n} \sigma_x & \text{for } n\text{-odd} \\ \frac{1}{2^n} I & \text{for } n\text{-even} \end{cases}$

Using these we can now evaluate:

$$\begin{aligned} R_{xy} &= e^{i g_{R_{xy}} \theta} = e^{i \frac{1}{2} \sigma_x \theta} = I + \frac{i}{2} \sigma_x \theta - \frac{1}{4} \frac{1}{2!} \sigma_x^2 \theta^2 + \frac{1}{3!} (-i) \frac{1}{8} \sigma_x \theta^3 + \dots \\ &= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \begin{pmatrix} i\theta & 0 \\ 0 & -i\theta \end{pmatrix} + \begin{pmatrix} -\frac{1}{4}\theta^2 & 0 \\ 0 & -\frac{1}{4}\theta^2 \end{pmatrix} + \begin{pmatrix} -\frac{i}{8}\theta^3 & 0 \\ 0 & \frac{i}{8}\theta^3 \end{pmatrix} + \dots \\ &= \begin{pmatrix} e^{i\theta} & 0 \\ 0 & e^{-i\theta} \end{pmatrix} \end{aligned}$$

4. To get an idea for how this works let's look at the most interesting case:

$$[g_4, g_5] = i f^{45k} g_k$$

$$\text{Left hand side: } [g_4, g_5] = \left[\frac{\lambda_4}{2}, \frac{\lambda_5}{2} \right] = \begin{pmatrix} 0 & 0 & i/2 \\ 0 & 0 & 0 \\ i/2 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & -i/2 \\ 0 & 0 & 0 \\ i/2 & 0 & 0 \end{pmatrix} - \begin{pmatrix} 0 & 0 & -i/2 \\ 0 & 0 & 0 \\ i/2 & 0 & 0 \end{pmatrix} \begin{pmatrix} 0 & 0 & i/2 \\ 0 & 0 & 0 \\ i/2 & 0 & 0 \end{pmatrix}$$

$$= \begin{pmatrix} i/4 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -i/4 \end{pmatrix} - \begin{pmatrix} -i/4 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & i/4 \end{pmatrix}$$

$$= \begin{pmatrix} i/2 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -i/2 \end{pmatrix}$$

$$\text{Right hand side: } i f^{45k} g_k = i f^{451} g_1 + i f^{452} g_2 + i f^{453} g_3 + i f^{454} g_4 + i f^{455} g_5 + i f^{456} g_6 + i f^{457} g_7 + i f^{458} g_8$$

$$= i f^{453} g_3 + i f^{458} g_8 \quad \text{but from what was given: } f^{345} = \frac{1}{2} \Rightarrow f^{453} = \frac{1}{2}$$

$$f^{458} = \frac{\sqrt{3}}{2} \quad \text{since 453 is a cyclic permutation of 345.}$$

$$= i \frac{1}{2} \lambda_3 + i \frac{\sqrt{3}}{2} \frac{1}{2} \lambda_8$$

$$= \frac{i}{4} \begin{pmatrix} 1 & & \\ & -1 & \\ & & 0 \end{pmatrix} + \frac{i\sqrt{3}}{4} \frac{1}{\sqrt{3}} \begin{pmatrix} 1 & & \\ & 1 & \\ & & -2 \end{pmatrix}$$

$$= \begin{pmatrix} i/4 & & \\ & -i/4 & \\ & & 0 \end{pmatrix} + \begin{pmatrix} i/4 & & \\ & i/4 & \\ & & -i/2 \end{pmatrix}$$

$$= \begin{pmatrix} i/2 & & \\ & 0 & \\ & & -i/2 \end{pmatrix} \quad \text{which agrees w/ l.h.s. } \checkmark$$

6. Recall that $\bar{J}_{\pm i} = \frac{1}{2}(\bar{J}_i \pm iK_i)$

We will demonstrate:

$$\begin{aligned} \text{a) } [\bar{J}_{+1}, \bar{J}_{+2}] &= \left[\frac{1}{2}(\bar{J}_1 + iK_1), \frac{1}{2}(\bar{J}_2 + iK_2) \right] = \frac{1}{4} \{ [\bar{J}_1, \bar{J}_2] - [K_1, K_2] + i[\bar{J}_1, K_2] + i[K_1, \bar{J}_2] \} \\ &= \frac{1}{4} \{ i\bar{J}_3 + i\bar{J}_3 - K_3 - K_3 \} & \text{Note: } [\bar{J}_i, K_j] &= i\epsilon^{ijk} K_k \Rightarrow [\bar{J}_1, K_2] = iK_3 \\ &= i\frac{1}{2}(\bar{J}_3 + iK_3) & [\bar{K}_i, \bar{J}_j] &= -[\bar{J}_j, \bar{K}_i] = -i\epsilon^{jik} K_k \\ &= i\bar{J}_{+3} & \rightarrow [\bar{K}_1, \bar{J}_2] &= -i\epsilon^{213} K_3 = iK_3 \end{aligned}$$

$$\begin{aligned} \text{b) } [\bar{J}_{+1}, \bar{J}_{-2}] &= \left[\frac{1}{2}(\bar{J}_1 + iK_1), \frac{1}{2}(\bar{J}_2 - iK_2) \right] = \frac{1}{4} \{ [\bar{J}_1, \bar{J}_2] + [K_1, K_2] - i[\bar{J}_1, K_2] + i[K_1, \bar{J}_2] \} \\ &= \frac{1}{4} \{ i\bar{J}_3 - i\bar{J}_3 + K_3 - K_3 \} \\ &= 0 \end{aligned}$$