

Review

TDSE: $i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \Psi}{\partial x^2} + V(x) \Psi$ (\hat{H}) - is ind. of t
 $\Psi(x,t)$ "stationary state"

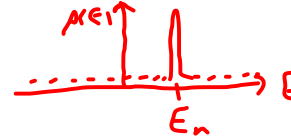
Assume $\Psi(x,t) = \sum_n c_n \phi_n(t) \psi_n(x)$

Then: TDSE $\Rightarrow \phi_n(t) = e^{-i \frac{E_n}{\hbar} t}$ (normalized)

$\Rightarrow \hat{H} \psi_n(x) = E_n \psi_n(x)$ TISE

$-\frac{\hbar^2}{2m} \frac{d^2}{dx^2} + V(x)$ Hamiltonian

$\rho(E) = \delta(E - E_n)$



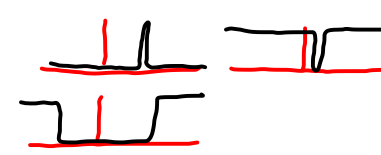
We will solve TISE for various $V(x)$:

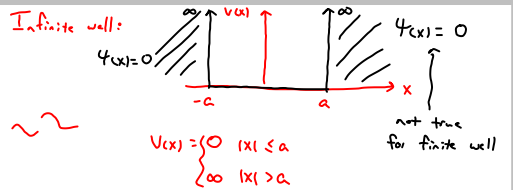
Infinite well  (simplest)

Harmonic osc.  (most important)

Zero (constant)  (non-normalizable ψ s)

Combine scattering & Bound states

$\left\{ \begin{array}{l} \delta\text{-function} \\ \text{Finite well} \end{array} \right.$ 



TISE $-\frac{\hbar^2}{2m} \frac{d^2 \psi_n(x)}{dx^2} + V(x) \psi_n(x) = E_n \psi_n(x)$

$|x| > a$ $V(x) = \infty$ $\frac{d^2 \psi}{dx^2} = \infty \psi$ HW 2.2
non-normalizable ($\psi = 0$)

$|x| < a$ $\frac{d^2 \psi_n(x)}{dx^2} = -\frac{2m E_n}{\hbar^2} \psi_n(x)$
 k^2

$\psi_n(x) = A \cos(kx) + B \sin(kx)$

Use special boundary conditions:

Continuity: $\psi(-a) = \psi(a) = 0$ matches $\psi(x)$ $|x| > a$

Smoothness: $\frac{\partial \psi}{\partial x}$ continuous \rightarrow allowed for singular potential

$\psi_n(a) = A \cos(ka) + B \sin(ka) = 0$

2 ways to solve: $A=0$ $k = \frac{n\pi}{2a}$ $n = \text{even}$ } Automatically
 $B=0$ $k = \frac{n\pi}{2a}$ $n = \text{odd}$ } satisfies $\psi(-a) = 0$

Normalizing:

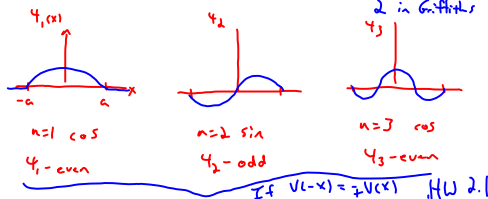
$\psi_n(x) = A \cos(kx) \Rightarrow A^2 \int_a^{-a} \cos^2(\frac{n\pi}{2a} x) dx = A^2 a = 1 \Rightarrow A = \frac{1}{\sqrt{a}}$

$\psi_n(x) = B \sin(kx) \Rightarrow B^2 \int_a^{-a} \sin^2(\frac{n\pi}{2a} x) dx = B^2 a = 1 \Rightarrow B = \frac{1}{\sqrt{a}}$

Solutions: $\psi_n(x) = \begin{cases} \frac{1}{\sqrt{a}} \cos(\frac{n\pi}{2a} x) & n = \text{odd} \\ \frac{1}{\sqrt{a}} \sin(\frac{n\pi}{2a} x) & n = \text{even} \end{cases}$

Energy levels (E_n)

Recall: $k^2 = \frac{2m E_n}{\hbar^2} \Rightarrow E_n = \frac{\hbar^2 k^2}{2m} = \frac{\hbar^2 \pi^2}{8ma^2} n^2$
2 in Griffiths



Finally TDSE:

$$\begin{aligned}\Psi(x,t) &= \sum_n c_n \phi_n(x) \psi_n(x) \\ &= \sum_{n=odd} c_n e^{-i \frac{E_n}{\hbar} t} \frac{1}{\sqrt{a}} \sin\left(\frac{n\pi}{2a} x\right) \\ &\quad + \sum_{n=even} c_n e^{-i \frac{E_n}{\hbar} t} \frac{1}{\sqrt{a}} \cos\left(\frac{n\pi}{2a} x\right)\end{aligned}$$

I promised: $\Psi(x,t) = \sum_n c_n \underbrace{\psi_n(x,t)}_{\text{orthonormal basis set}}$

$$\psi_n(x,t) = \phi_n(x) \psi_n(x)$$

Orthogonality: $\int_{-a}^a \psi_m^*(x,t) \psi_n(x,t) dx \stackrel{?}{=} 0 \quad n \neq m$

$$\int_{-a}^a \psi_m^*(x) \psi_n(x) dx$$

$$\text{Use: } \int_{-a}^a \sin\left(\frac{n\pi}{2a} x\right) \sin\left(\frac{m\pi}{2a} x\right) dx = 0 \quad \text{if } n \neq m$$

$$\int_{-a}^a \sin\left(\frac{n\pi}{2a} x\right) \cos\left(\frac{m\pi}{2a} x\right) dx = 0 \quad \text{trivially}$$

$$\int_{-a}^a \cos\left(\frac{n\pi}{2a} x\right) \cos\left(\frac{m\pi}{2a} x\right) dx = 0 \quad \text{if } n \neq m$$

$$\int_{-a}^a \psi_m^*(x) \psi_n(x) dx = 0 \quad n \neq m$$

Normality $\int_{-a}^a \psi_n^*(x,t) \psi_n(x,t) dx \stackrel{?}{=} 1$

$$\int_{-a}^a \psi_n^2(x) dx = 1$$

Generally: $\int_{-a}^a \psi_m^*(x,t) \psi_n(x,t) dx = \delta_{mn} = \begin{cases} 1 & n=m \\ 0 & n \neq m \end{cases}$

$$\text{For } \Psi(x,t): \int_{-\infty}^{\infty} \Psi^*(x,t) \Psi(x,t) dx = c_1$$

↑
overlap of Ψ w/ ψ_1

What are c_n ? $|c_n|^2$ = probability of measuring E_n for $\Psi(x,t)$

$$\begin{aligned}\text{Note: } \int_{-a}^a \Psi^*(x,0) \Psi(x,0) dx &= \int_{-a}^a \left(\sum_n c_n \psi_n(x) \right) \left(\sum_m c_m^* \psi_m^*(x) \right) dx \\ &= \sum_n \sum_m c_n^* c_m \underbrace{\int_{-a}^a \psi_n^*(x) \psi_m(x) dx}_{\delta_{nm}} \\ &= \sum_n |c_n|^2 = 1 \quad \left\langle \hat{H} \right\rangle = E_{avg} = \sum_n |c_n|^2 E_n\end{aligned}$$