

Review: Classical SHO $H = \frac{p^2}{2m} + \frac{1}{2}m\omega^2 x^2$
 $\dot{x} = \frac{\partial H}{\partial p}$ $\dot{p} = -\frac{\partial H}{\partial x}$

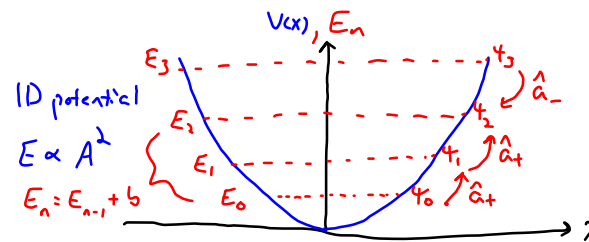
Quantum SHO $\hat{H} = \frac{\hat{p}^2}{2m} + \frac{1}{2}m\omega^2 \hat{x}^2$
 TISE $\hat{H}\psi_n = E_n \psi_n$

$\hat{H} = \hbar\omega(\hat{a}_- \hat{a}_+ + \frac{1}{2})$
 $[\hat{x}, \hat{p}] = i\hbar$ $\hat{a}_- \hat{a}_+ = \hat{a}_+ \hat{a}_- + 1$
 $\hat{a}_\pm = \frac{1}{\sqrt{2\hbar m\omega}}(m\omega \hat{x} \mp i\hat{p}) \Rightarrow [\hat{a}_-, \hat{a}_+] = 1$

$\hat{H} = \hbar\omega(\hat{a}_+ \hat{a}_- + \frac{1}{2})$

\Downarrow

$\hat{H}\psi_n = E_n \psi_n$



Our job is to find normalized ψ_n and E_n .

\hat{a}_+ - raising } ladder operators (creation
 \hat{a}_- - lowering } annihilation

useful in angular momentum
 and in QFT

Claim: All we need to do is find ψ_0 and E_0
 Everything else will come from the algebra
 of \hat{a}_+, \hat{a}_- .

Start by proving the "stepping":

$$\text{Assume: } \hat{H} \psi_n = E_n \psi_n$$

$$\text{Consider: } \hat{H}(\hat{a}_+ \psi_n) = E_{n+1}(\hat{a}_+ \psi_n)$$

$$\begin{aligned} (a_+ a_- + \frac{1}{2}) a_+ &= \hbar \omega (\hat{a}_+ \hat{a}_- + \frac{1}{2}) \hat{a}_+ \psi_n \\ a_+ a_- a_+ + \frac{1}{2} a_+ & \\ a_+ (a_- a_+ + \frac{1}{2}) &= \hat{a}_+ \hbar \omega (\hat{a}_- \hat{a}_+ + \frac{1}{2}) \psi_n \\ & \quad \hat{a}_+ \hat{a}_- + 1 \\ &= \hat{a}_+ \hbar \omega (\hat{a}_+ \hat{a}_- + \frac{1}{2} + 1) \psi_n \\ & \quad \hat{H} \\ \psi_{n+1} &= \hat{a}_+ (E_n + \hbar \omega) \psi_n \\ \hat{H}(\hat{a}_+ \psi_n) &= (E_n + \hbar \omega) (\hat{a}_+ \psi_n) \\ & \quad \text{is a solution to TISE} \end{aligned}$$

$$\text{You derive: } \hat{H}(\hat{a}_- \psi_n) = (E_n - \hbar \omega) (\hat{a}_- \psi_n)$$

$$\hat{H} \psi_{n+1} = E_{n+1} \psi_{n+1} \quad \psi_{n+1} \propto \hat{a}_+ \psi_n$$
$$\psi_{n-1} \propto \hat{a}_- \psi_n$$

Caution: $\psi_{n+1} = \hat{a}_+ \psi_n$ is a solution to TISE
 but not nec. normalized

$$\begin{aligned}
 -i\hbar \int_{-\infty}^{\infty} \psi_{n+1}^* \psi_{n+1} dx &= 1 \\
 -i\hbar \frac{d}{dx} & \\
 &= \int_{-\infty}^{\infty} (\hat{a}_+ \psi_n)^* (\hat{a}_+ \psi_n) dx \\
 &= \int_{-\infty}^{\infty} \frac{1}{\sqrt{2m\hbar\omega}} (m\omega\hat{x} - i\hat{p}) \psi_n^* \hat{a}_+ \psi_n dx \\
 &= \int_{-\infty}^{\infty} \frac{1}{\sqrt{2m\hbar\omega}} (m\omega\hat{x} + \hbar \frac{d}{dx}) \psi_n^* \hat{a}_+ \psi_n dx
 \end{aligned}$$

We know $\hat{x}\psi = \psi x$

But $\frac{d}{dx}$ is tricky:

$$\begin{aligned}
 &\int_{-\infty}^{\infty} (\frac{d}{dx} \psi) (c \psi) dx \\
 &\frac{d}{dx} (\psi c \psi) - \psi \frac{d(c \psi)}{dx} \\
 &\int_{-\infty}^{\infty} \frac{d}{dx} (\psi c \psi) dx - \int_{-\infty}^{\infty} \psi \frac{d(c \psi)}{dx} dx \\
 &\psi c \psi \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} \psi \frac{d(c \psi)}{dx} dx
 \end{aligned}$$

$$= \int_{-\infty}^{\infty} \psi_n^* \frac{1}{\sqrt{2m\hbar\omega}} (m\omega\hat{x} + \hbar \frac{d}{dx}) \hat{a}_+ \psi_n dx$$

$$= \int_{-\infty}^{\infty} \psi_n^* \underbrace{\hat{a}_- \hat{a}_+}_{\frac{\hat{H}}{m\omega} + \frac{1}{2}} \psi_n dx$$

$$= \int_{-\infty}^{\infty} \psi_n^* \left(\frac{\hat{H}}{m\omega} + \frac{1}{2} \right) \psi_n dx$$

$$= \int_{-\infty}^{\infty} \psi_n^* \left(\frac{E_n}{m\omega} + \frac{1}{2} \right) \psi_n dx$$

$$= \left(\frac{E_n}{m\omega} + \frac{1}{2} \right) \int_{-\infty}^{\infty} \psi_n^* \psi_n dx$$

$$\int_{-\infty}^{\infty} (\hat{a}_+ \psi_n)^* (\hat{a}_+ \psi_n) dx = \frac{E_n}{m\omega} + \frac{1}{2}$$

$$\psi_{n+1} = \frac{1}{\sqrt{\frac{E_n}{m\omega} + \frac{1}{2}}} \hat{a}_+ \psi_n$$

\hat{L} gives a normalized solution to TISE iff ψ_n is a norm. sol. to TISE

You prove:
$$\psi_{n-1} = \frac{1}{\sqrt{\frac{E_n}{m\omega} - \frac{1}{2}}} \hat{a}_- \psi_n$$

If I had a normalized ψ_0 and I knew E_0 ,
then I can build all $\psi_{n>0}$ and know $E_n!!!$

$$\psi_1 = \psi_{0+1} = \frac{1}{\sqrt{\frac{E_0}{\hbar\omega} + \frac{1}{2}}} \hat{a}_+ \psi_0 \quad E_1 = E_0 + \hbar\omega$$

$$\psi_2 = \psi_{1+1} = \frac{1}{\sqrt{\frac{E_1}{\hbar\omega} + \frac{1}{2}}} \hat{a}_+ \psi_1$$

$$= \frac{1}{\sqrt{\frac{E_1}{\hbar\omega} + \frac{1}{2}}} \hat{a}_+ \frac{1}{\sqrt{\frac{E_0}{\hbar\omega} + \frac{1}{2}}} \hat{a}_+ \psi_0$$

$$\psi_0, E_0 = ?$$

Recall from HW 2.2 if $E_n < V_{min}$ then ψ_n is BADD!!!

Assume $E_n > V_{min} \stackrel{SHO}{=} 0$

But $\hat{a}_- \psi_n \Rightarrow E_n - \hbar\omega = E_{n-1}$

This process has to bottom out!!

Impose $\hat{a}_- \psi_0 = 0 \Rightarrow \psi_{-1} = 0$
 $\psi_{-2} = \hat{a}_- \psi_{-1} = 0$

Equation can be solved for ψ_0

$$\frac{1}{\sqrt{2\pi}\hbar} (m\hat{x} + i\hat{p})\psi_0 = 0$$

$$\frac{1}{\sqrt{2\pi}\hbar} (m\hat{x} + \hbar \frac{d}{dx})\psi_0 = 0 \Rightarrow \frac{d\psi_0}{dx} = -\frac{m\omega}{\hbar} \hat{x} \psi_0$$

$$\psi_0(x) = A e^{-\frac{m\omega}{2\hbar} x^2}$$

$$\int \psi_0^* \psi_0 dx = 1$$

$$\Downarrow$$

$$A = \left(\frac{m\omega}{\pi\hbar}\right)^{\frac{1}{4}}$$

Finding E_0 : $\hat{H}\psi_0 = E_0\psi_0$

$$\hbar\omega(\hat{a}_+ \hat{a}_- + \frac{1}{2})\psi_0 = E_0\psi_0$$

$$\frac{1}{2}\hbar\omega\psi_0 = E_0\psi_0$$

$$E_0 = \frac{1}{2}\hbar\omega$$

Back to normalization of $\psi_{n+1} = \frac{1}{\sqrt{E_{n+1} - E_n}} \hat{a}_+ \psi_n$

$$\psi_1 = \frac{1}{\sqrt{E_1 - E_0}} \hat{a}_+ \psi_0$$

$$\psi_1 = \hat{a}_+ \psi_0$$

$$\psi_2 = \frac{1}{\sqrt{E_2 - E_1}} \hat{a}_+ \psi_1$$

$$= \frac{1}{\sqrt{2}} \hat{a}_+ \psi_1$$

$$\psi_2 = \frac{1}{\sqrt{2!}} \hat{a}_+ \psi_1$$

$$= \frac{1}{\sqrt{n!}} \hat{a}_+ \psi_0$$

Why did this work?

$$-\frac{\hbar^2}{2m} \frac{\partial^2 \psi_n}{\partial x^2} + \frac{1}{2} m \omega^2 x^2 \psi_n = E_n \psi_n$$

This sucks!!

We "factorized" $\hat{H} = \hbar\omega(\hat{a}_+ \hat{a}_- + \frac{1}{2})$

$$\hat{H} \psi_0 = E_0 \psi_0$$

$$\hbar\omega(\hat{a}_+ \hat{a}_- + \frac{1}{2}) \psi_0 = E_0 \psi_0$$

$$\hat{a}_- \psi_0 = 0$$