

Aug. 75.02

High 92.5

In stationary states $\frac{d\langle \hat{p} \rangle}{dt} = 0$ $\psi(x,t) = e^{iEt/\hbar}$ $\psi^*(x,t) = e^{-iEt/\hbar}$

$$\Psi = \sum_n c_n \psi_n \approx c_1 \psi_1 + c_2 \psi_2$$

$$\langle \hat{x} \rangle = \int \Psi^* x \Psi dx = \int (c_1^* \psi_1^* + c_2^* \psi_2^*) x (c_1 \psi_1 + c_2 \psi_2) dx$$

$$\sigma_x \sigma_p \geq \frac{\hbar}{2}$$

$$\sigma_x = \sqrt{\langle \hat{x}^2 \rangle - \langle \hat{x} \rangle^2}$$

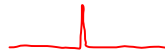
$$\sigma_p = \sqrt{\langle \hat{p}^2 \rangle - \langle \hat{p} \rangle^2}$$

$$\langle \hat{A} \rangle = \int \psi^* \hat{A} \psi dx \quad \langle \hat{A}^2 \rangle \quad \sigma_A = \sqrt{\langle \hat{A}^2 \rangle - \langle \hat{A} \rangle^2} = 0$$

$$\int E \cdot 0 dx$$

$$0 \cdot 0$$

$$E \cdot E$$



$$\psi(x) \text{ and } V(x) \Rightarrow \hat{H} \psi(x) = E \psi(x)$$
$$\left(-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + V \right) \psi(x) = E \psi(x)$$

$$[\hat{x}, \hat{p}^2] \psi(x) = 2\hbar^2 \frac{\partial}{\partial x} \psi(x)$$

$$[\hat{x}, \hat{p}] = i\hbar \quad \hat{a}_+ = \hat{p} + i\hat{x}$$
$$\hat{a}_- = \hat{p} - i\hat{x}$$

$$\langle \hat{x} \hat{p} \rangle \quad \langle \hat{x}^2 \rangle, \langle \hat{p}^2 \rangle, \langle \hat{x}^4 \rangle, \langle \hat{p}^4 \rangle$$



$$Ae^{ikx} + Be^{-ikx}$$
$$A \sin(kx) + B \cos(kx)$$

continuous spectrum

$$\phi(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \psi(x, t) \Psi^*(x, 0) dx$$

$$\text{TISE} \Rightarrow \psi_n(x) \quad \psi(k, x)$$

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apply b.c.s

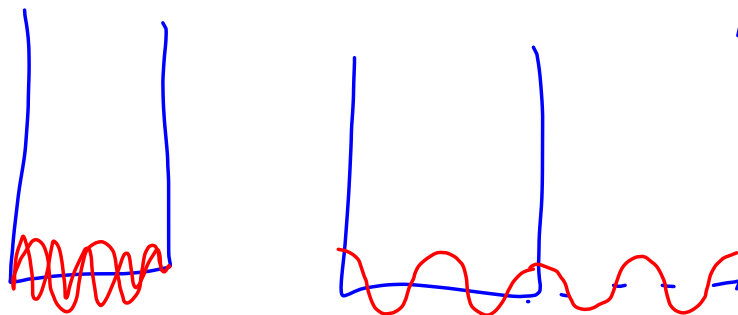
$$\begin{aligned} \sin(kx) \\ \cos(kx) \end{aligned} \Rightarrow k = \frac{2n\pi}{L}$$

$$\psi_n(x) = A \cos\left(\frac{2n\pi}{L} x\right)$$

$$\psi(x, 0) = A \cos\left(\frac{6\pi}{L} x\right)$$

$$C_3 = 1 \quad C_{n \neq 3} = 0$$

$$V(x) = \frac{1}{2} \pi \omega^2 x^2$$



$$\hat{a}_+ = \frac{1}{\sqrt{2\pi\hbar\omega}} (\hat{x}m\omega - i\hat{p}) \quad [\hat{x}, \hat{p}]$$

$$\hat{a}_- = \frac{1}{\sqrt{2\pi\hbar\omega}} (\hat{x}m\omega + i\hat{p})$$

$$S = \int L dt \quad L = T - V = \left[ \frac{1}{2} m \dot{x}^2 - V(x) \right] dt$$

$$\delta S = \delta \int L dt = 0 \quad x(t)$$

$$f(x) = x^3 \quad \dot{x}(t) = v(t)$$

$$g(f(x))$$



e.o.m.

Lagrange e.o.m.

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}} \right) - \frac{\partial L}{\partial x} = 0$$

$$PI = \int e^{\frac{i}{\hbar} \int L dt} \rightarrow \text{T.D.S.E.}$$

$$\text{EM wave} \quad \frac{\partial^2 \epsilon}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2 \epsilon}{\partial t^2} = 0 \Rightarrow \epsilon(x,t) = E_0 e^{i(kx - \omega t)}$$

$$(-k^2 + \frac{\omega^2}{c^2}) \epsilon(x,t) = 0 \quad \frac{\partial^2 \epsilon}{\partial x^2} = -k^2 \epsilon(x,t)$$

$$\frac{\partial^2 \epsilon}{\partial t^2} = -\omega^2 \epsilon(x,t)$$

$$k^2 = \frac{\omega^2}{c^2} \Rightarrow c = \frac{\omega}{k} = \lambda f$$

Consider:  $\epsilon = \hbar \omega$ ,  $p = \hbar k$

$$\left[ -\frac{\hbar^2}{2m} - \frac{1}{2} \left( -\frac{\hbar^2}{2m} \right) \right] \epsilon(x,t) = 0$$

$$-\frac{\hbar^2}{2m} \left[ p^2 - \frac{\epsilon^2}{c^2} \right] \epsilon(x,t) = 0$$

$$\epsilon^2 = p^2 c^2$$

↓ massive

$$\epsilon^2 = p^2 c^2 + m^2 c^4$$

Apply de Broglie  $E_0 \rightarrow \phi_0 \quad \frac{i}{\hbar} (px - \epsilon t)$

$$-\frac{\hbar^2}{2m} \left[ p^2 - \frac{\epsilon^2}{c^2} + m^2 c^2 \right] \phi_0 e^{\frac{i}{\hbar} (px - \epsilon t)} = 0$$

$$\left( \frac{\partial^2}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \frac{\hbar^2 c^2}{2m} \right) \phi_0 e^{\frac{i}{\hbar} (px - \epsilon t)} = 0$$

$$\frac{\partial^2 \phi}{\partial x^2} - \frac{\hbar^2 c^2}{2m} \phi = \frac{1}{c^2} \frac{\partial^2 \phi}{\partial t^2} \quad \left[ \text{Klein-Gordon eqn.} \right]$$

relativistic SE

$$\text{NR limit: } \epsilon = \hbar c^2 \sqrt{1 + \frac{p^2}{\hbar^2 c^2}} \approx \hbar c^2 \left( 1 + \frac{1}{2} \frac{p^2}{\hbar^2 c^2} \right)$$

$$= \hbar c^2 + \frac{1}{2} \frac{p^2}{m}$$

$$= \hbar c^2 + T$$

$$\begin{aligned} \phi(x,t) &= \phi_0 e^{\frac{i}{\hbar}(\rho x - \epsilon t)} = \phi_0 e^{\frac{i}{\hbar}(\rho x - \hbar c t - T t)} \\ &= e^{-\frac{i \hbar c t}{\hbar}} e^{\frac{i}{\hbar}(\rho x - T t)} \end{aligned}$$

oscillates for relativistic system

Introduce:  $\phi = e^{-\frac{i}{\hbar} \hbar c t} \psi(x,t)$

$$\begin{aligned} \frac{\partial^2 \phi}{\partial t^2} &= -\frac{1}{\hbar^2} \hbar^2 c^2 e^{-\frac{i}{\hbar} \hbar c t} \psi \\ &\quad - \frac{2i}{\hbar} \hbar c e^{-\frac{i}{\hbar} \hbar c t} \frac{\partial \psi}{\partial t} + e^{-\frac{i}{\hbar} \hbar c t} \frac{\partial^2 \psi}{\partial t^2} \end{aligned}$$

$$\left( \frac{\partial^2}{\partial x^2} - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \frac{\hbar^2 c^2}{\hbar^2} \right) \phi = 0$$

$$e^{-\frac{i}{\hbar} \hbar c t} \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{2i \hbar}{\hbar} \frac{\partial \psi}{\partial t} \right) = 0$$

TDSE