

Review: Spherical Harmonics  $Y_l^m(\theta, \phi)$  (angular solutions to TISE for  $V(r)$  incl.  $V_{hyd.}$ )

$Y_l^m$  are eigenfunctions of  $\hat{L}^2$  and  $\hat{L}_z$  — are hermitian observables

$$\begin{aligned} \hat{L}^2 Y_l^m &= \hbar^2 l(l+1) Y_l^m \\ \hat{L}_z Y_l^m &= \hbar m Y_l^m \end{aligned}$$

$l = 0, 1, 2, \dots$   
 $m = -l, \dots, +l$

explains why these are orthogonal

$\vec{L} = \sum_{i=1}^3 \vec{r}_i \times \vec{p}_i$  (motion of CM)


Classically we can split orbital motion from spinning motion.

$\vec{S} = I_{cm} \vec{\omega}$  (about the CM)

The  $Y_l^m$  describe only the orbital motion of electrons in atoms.

Q. Can the electron spin?

Classically  ~~$\vec{L} = \sum \vec{r}_i \times \vec{p}_i$~~   $\vec{L} = 0$   $\sim \infty$  for a nonzero  $\vec{S}$  doesn't work for spin



Spin is an intrinsic and unchangeable property of point particles.

nothing changes it!  
this is not built from external or spacetime coord. (no built from  $x, p$ )

We have an abstract algebraic approach to angular momentum.

$$\begin{aligned} \hat{S}^2 \chi_s^m &= \hbar^2 s(s+1) \chi_s^m & s &= 0, \frac{1}{2}, 1, \frac{3}{2}, \dots \\ \hat{S}_z \chi_s^m &= \hbar m \chi_s^m & m &= -s, \dots, s \end{aligned}$$

For  $\vec{L} = \frac{\hbar}{i} \left[ r(\partial_x \partial_\phi + \partial_\phi \partial_x) + (\hat{r} \times \hat{\theta}) \frac{\partial}{\partial \theta} + (\hat{r} \times \hat{\phi}) \frac{1}{\sin \theta} \frac{\partial}{\partial \phi} \right]$   
from  $\vec{r} \times \vec{p}$

$\Rightarrow Y_l^m(\theta, \phi)$  labelled by  $m$  and  $l$ .

For spin we only have labels  $m$  and  $s$ !!

We would like to know  $\chi_s^m$  and  $\hat{S}^2$  and  $\hat{S}_z$   
 starting from eigenvalue expressions with known eigenvalues.

Context:  $\vec{L}$  where we know  $\hat{L}^2, \hat{L}_z$  then found  
 eigenfunctions and eigenvalues.

Example:  $s = \frac{1}{2} \Rightarrow m = -\frac{1}{2}, \frac{1}{2}$   
 $\uparrow$  spin down  
 $\uparrow$  spin up

$$\chi_{\frac{1}{2}}^{-\frac{1}{2}} = |\frac{1}{2}, -\frac{1}{2}\rangle = \chi_- = \begin{pmatrix} 0 \\ 1 \end{pmatrix} \quad \text{Spin down} \rightarrow$$

$$\chi_{\frac{1}{2}}^{\frac{1}{2}} = |\frac{1}{2}, \frac{1}{2}\rangle = \chi_+ = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad \text{Spin up} \rightarrow$$

Building  $\hat{S}^2$  and  $\hat{S}_z$ :

a)  $\hat{S}^2 \chi_+ = \hbar^2 \frac{1}{2} (\frac{1}{2} + 1) \chi_+ = \frac{3}{4} \hbar^2 \chi_+$

b)  $\hat{S}^2 \chi_- = \hbar^2 \frac{1}{2} (\frac{1}{2} + 1) \chi_- = \frac{3}{4} \hbar^2 \chi_-$

Suppose  $\hat{S}^2 = \begin{pmatrix} a & d \\ 0 & f \end{pmatrix}$   $\leftarrow$  complex

a)  $\begin{pmatrix} a & d \\ 0 & f \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \frac{3}{4} \hbar^2 \begin{pmatrix} 1 \\ 0 \end{pmatrix} \Rightarrow \begin{pmatrix} a \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{3}{4} \hbar^2 \\ 0 \end{pmatrix}$

b)  $\begin{pmatrix} a & d \\ 0 & f \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \frac{3}{4} \hbar^2 \begin{pmatrix} 0 \\ 1 \end{pmatrix} \Rightarrow \begin{pmatrix} d \\ f \end{pmatrix} = \begin{pmatrix} 0 \\ \frac{3}{4} \hbar^2 \end{pmatrix}$

$$\Rightarrow \hat{S}^2 = \begin{pmatrix} \frac{3}{4} \hbar^2 & 0 \\ 0 & \frac{3}{4} \hbar^2 \end{pmatrix} = \frac{3}{4} \hbar^2 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

Now for  $\hat{S}_z$ : let  $\hat{S}_z = \begin{pmatrix} g & r \\ 0 & t \end{pmatrix}$

$\hat{S}_z \chi_+ = \frac{\hbar}{2} \chi_+ \Rightarrow \begin{pmatrix} g & r \\ 0 & t \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} \frac{\hbar}{2} \\ 0 \end{pmatrix} = \begin{pmatrix} g \\ 0 \end{pmatrix}$

$\hat{S}_z \chi_- = -\frac{\hbar}{2} \chi_- \Rightarrow \begin{pmatrix} g & r \\ 0 & t \end{pmatrix} \begin{pmatrix} 0 \\ 1 \end{pmatrix} = \begin{pmatrix} 0 \\ -\frac{\hbar}{2} \end{pmatrix} = \begin{pmatrix} 0 \\ t \end{pmatrix}$

$$\Rightarrow \hat{S}_z = \begin{pmatrix} \frac{\hbar}{2} & 0 \\ 0 & -\frac{\hbar}{2} \end{pmatrix} = \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

2 issues remain

i) With  $\vec{L}$  we could also build  $\hat{L}_x, \hat{L}_y$  so can we build  
 $\hat{S}_x, \hat{S}_y$ ?

ii) The end quantumness of  $\vec{S}$  comes in  $[\hat{S}_x, \hat{S}_y] = i \hbar \hat{S}_z$ .

So building  $\hat{S}_x$  and  $\hat{S}_y$ :

Do not have an eigenvalue eqn.  ~~$\hat{S}_x \chi_{\pm} = \lambda \chi_{\pm}$~~

Note:  $\hat{S}_x$  will have its own eigenstates  $\chi_{\pm}^x \neq \chi_{\pm}$ .

Recall:  $\hat{S}_{\pm} = \hat{S}_x \pm i \hat{S}_y \Rightarrow \begin{cases} \hat{S}_x = \frac{1}{2} (\hat{S}_+ + \hat{S}_-) \\ \hat{S}_y = \frac{i}{2} (\hat{S}_+ - \hat{S}_-) \end{cases}$

from  $\hat{S}_z$   $\leftarrow$  from HW  
 We know:  $\begin{cases} \hat{S}_z \chi_{\pm} = \pm \frac{\hbar}{2} \chi_{\pm} \\ \hat{S}_+ \chi_- = \hbar \chi_+ \\ \hat{S}_- \chi_+ = \hbar \chi_- \end{cases} \Rightarrow \begin{cases} \hat{S}_+ = \begin{pmatrix} 0 & \hbar \\ 0 & 0 \end{pmatrix} \\ \hat{S}_- = \begin{pmatrix} 0 & 0 \\ \hbar & 0 \end{pmatrix} \end{cases}$

$$\Rightarrow \hat{S}_x = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$\hat{S}_y = \frac{\hbar}{2} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$$

$$\text{Finally: } \hat{S} = \frac{\hbar}{2} \hat{\sigma}$$

$$\hat{\sigma} = (\sigma_x, \sigma_y, \sigma_z)$$

$$= \left( \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right)$$

Pauli spin matrices

3 things to consider

$$i) \text{ Does } [\hat{S}_x, \hat{S}_y] = i\hbar \hat{S}_z ?$$

ii) Are eigenspinors of  $\hat{S}_x$  ( $\chi_{\pm}^x$ ) different from  $\chi_{\pm}$  of  $\hat{S}_z$ ?

$$iii) \text{ Is } \hat{S}^2 \chi_{\pm}^x = \frac{3\hbar^2}{4} \chi_{\pm}^x$$

$$\begin{aligned} i) \frac{\hbar^2}{4} \left( \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} - \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \right) &= \frac{\hbar^2}{4} \left( \begin{pmatrix} i & 0 \\ 0 & -i \end{pmatrix} - \begin{pmatrix} -i & 0 \\ 0 & i \end{pmatrix} \right) \\ &= \frac{\hbar^2}{4} \begin{pmatrix} 2i & 0 \\ 0 & -2i \end{pmatrix} \\ &= i\hbar \frac{\hbar}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ &= i\hbar \hat{S}_z \end{aligned}$$

$$\therefore \hat{S}_x = \frac{\hbar}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \Rightarrow \chi_{\pm}^x = \begin{pmatrix} 1/\sqrt{2} \\ \pm 1/\sqrt{2} \end{pmatrix} \neq \chi_{\pm}$$

$$iii) \hat{S}^2 \chi_{\pm}^x = \frac{3}{4} \hbar^2 \underbrace{\begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}}_{\mathbb{I}} \chi_{\pm}^x = \frac{3}{4} \hbar^2 \chi_{\pm}^x$$

We have a spin- $\frac{1}{2}$  particle  $\chi = a\chi_+ + b\chi_-$

$$\langle \hat{S}_z \rangle_{\chi} = \chi^\dagger \hat{S}_z \chi$$

= ( More Next time ...