

Reminder: HW Quiz this Wednesday

$l = 0, \dots, n-1$
 $m_l = -l, \dots, l$ $\psi_{n, m_l, l}$ are "hydrogenic" wavefunctions ($e \rightarrow Ze$)
 Periodic Table
 $Z=1$ Hydrogen $\psi_{100} \chi_{\pm}$
 $Z=2$ Helium $\psi_{100} \psi_{100} \frac{1}{\sqrt{2}} (\chi_+^1 \chi_-^1 - \chi_-^1 \chi_+^1)$
 $Z=3$ Lithium $\psi_{100} \psi_{100} \psi_{200} \underbrace{\chi_+^1 \chi_+^1 \chi_+^1}_{\text{with symmetric must be antisymmetric}}$

Larger l values (for a given n) are higher in energy due to inner-electron screening of nucleus.

$Z=3 \Rightarrow \psi_{100} \psi_{100} \psi_{200}$

For $Z=19$ (potassium) the $n=3, l=2$ state is higher in energy than $n=4, l=0$!

Nonrelativistic: $l = 0, 1, 2, \dots$
 $l = 0 (s), 1 (p), 2 (d), 3 (f), 4 (g), \dots$

$(1s)^2 (2s)^2 (2p)^2$ carbon $Z=6$
 $(1s)^2 (2s)^2 (2p)^3$
 $L=0 \rightarrow l=1 \Rightarrow m_l = -1, 0, 1 \Rightarrow$ can hold 6 total electrons
 $\underbrace{\psi_{100}^2 \psi_{200}^2 \frac{1}{\sqrt{2}} (\chi_+^1 \chi_-^1 - \chi_-^1 \chi_+^1)}_{S=0}$
 \downarrow
 $(1s)^2 (2s)^2 (2p)^2 \xrightarrow{L=1}$
 \downarrow
 $\psi_{200}^3 \psi_{200}^1 \frac{1}{\sqrt{2}} (\chi_+^1 \chi_-^1 - \chi_-^1 \chi_+^1) \xrightarrow{S=0}$

Total cons. num. $l_1=0, l_2=0, s_1+s_2=0$ } All together
 $3, 1 \quad l_3=0, l_4=0, s_3+s_4=0$ } + l_5, s_5
 $5, 6 \quad l_5=1, l_6=1$ } 0 cons.

required sym: $\psi_{210}^5 \psi_{210}^6 \psi_{211}^5 \psi_{211}^6 \Rightarrow s_1+s_2=0$
 sym. or antisym: $\psi_{210}^5 \psi_{211}^6 \Rightarrow s_1+s_2=0, 1$

When adding $L^5 + L^6 = 2, 1, 0$ w/ $S^5 + S^6 = 0, 1$

then $J_{tot} = 3, 2, 1, 0$
 $\underbrace{2S+1}_{\text{total spin}}$ Hund's Rules (problem 5.13 in book)
 $\underbrace{L}_{\text{total ang. num.}}$
 $\underbrace{J}_{\text{total orb. cons. num.}}$

Solids

Focus is on detached loosely bound valence electrons which roam the solid.

2 approximations

Free-electron gas: Treats electrons as P.I.B.

Block's theory: Periodic potential



Free-electron Gas

$$V(x,y,z) = \begin{cases} 0 & 0 < x,y,z < L_x, L_y, L_z \\ \infty & \text{otherwise} \end{cases}$$

From H.W. # 4.2

$$\psi_{n_x, n_y, n_z} = \sqrt{\frac{8}{L_x L_y L_z}} \sin\left(\frac{n_x \pi}{L_x} x\right) \sin\left(\frac{n_y \pi}{L_y} y\right) \sin\left(\frac{n_z \pi}{L_z} z\right)$$
$$n_x, n_y, n_z = 1, 2, 3, \dots$$

Note: $\psi(0) = 0 = \psi(x=L_x, y, z) = \psi(x, y=L_y, z)$

$$E_{n_x, n_y, n_z} = \frac{\hbar^2}{2m} \left(\frac{n_x^2 \pi^2}{L_x^2} + \frac{n_y^2 \pi^2}{L_y^2} + \frac{n_z^2 \pi^2}{L_z^2} \right) = \frac{\hbar^2}{2m} k^2$$
$$\vec{k} = \left(\frac{n_x \pi}{L_x}, \frac{n_y \pi}{L_y}, \frac{n_z \pi}{L_z} \right)$$

Let's go to k -space (also called reciprocal space)

N - # of atoms in solid (big for macroscopic)

q - # of free-electrons per atom (typically ~ 2)

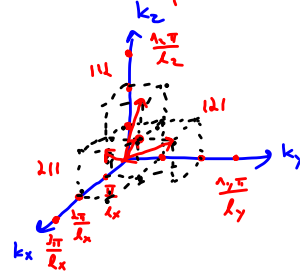
Nq - total # of free electrons in sample

$$\rho = \frac{Nq}{V} = \frac{Nq}{L_x L_y L_z} = \text{free-electron density}$$

Electrons are fermions \Rightarrow Pauli exclusion principle

\Rightarrow 2 electrons in each ψ_{n_x, n_y, n_z} state.

k-space is the space of allowed \vec{k} -vectors



At low T

$$N_q = 1 \quad \hat{k} = (1, 1, 1)$$

$$N_q = 2 \quad \hat{k} = (1, 1, 1), (1, 1, 2)$$

$$N_q = 3 \quad \hat{k}_{1,2} = (1, 1, 1), \hat{k}_3 = (1, 1, 2), (1, 2, 1), (2, 1, 1)$$

⋮

$$N_q = 8 \quad \text{can live in } (1, 1, 1), (1, 1, 2), (1, 2, 1), (2, 1, 1)$$

As $N_q \rightarrow \text{large}$ we fill up an octant of k-space.

