

Fermi-electron gas in 1D

$$\psi_n = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi}{L} x\right)$$

$$E_n = \frac{\hbar^2}{2m} \frac{n^2 \pi^2}{L^2} = \frac{\hbar^2}{2m} k^2 \quad \vec{k} = \frac{n\pi}{L} \hat{x}$$

For $\frac{N_F}{2}$ total pairs we populate $\frac{N_F}{2} \frac{\pi}{L}$ volume of k-space at $T \rightarrow 0$.

$$\text{Then: } k_{\text{max}} = \frac{N_F}{2} \frac{\pi}{L} = k_F \Rightarrow E_F = \frac{\hbar^2}{2m} \frac{N_F^2 \pi^2}{4L^2}$$

$$E_{\text{tot}} = \int dE = \int E dN = \int_0^{k_F} \underbrace{E}_{\frac{\hbar^2 k^2}{2m}} \underbrace{dV}_{\frac{L}{\pi} dk} = \int_0^{k_F} \frac{\hbar^2}{2m} k^2 \frac{L}{\pi} dk$$

$$= \frac{L \hbar^2}{3m\pi} k_F^3$$

$$= \frac{\hbar^2}{24m} \frac{1}{L^2} N_F^3 \Rightarrow P = -\frac{dE_{\text{tot}}}{dL}$$

For bosons

$$E_{\text{tot}} = N \frac{\hbar^2 \pi^2}{2m L^2}$$

$$\begin{aligned} \text{1D: } E_{\text{tot}}^b &\sim N_b & E_{\text{tot}}^f &\sim N_f^3 \\ \text{3D: } E_{\text{tot}}^b &\sim N_b & E_{\text{tot}}^f &\sim N_f^{5/3} \end{aligned}$$

In higher dimensions the distribution of states will "spread" out in k-space.



$$E_{\text{tot}} \propto k^2 \Rightarrow E_{\text{tot}}^{\frac{1}{2}} \propto k$$

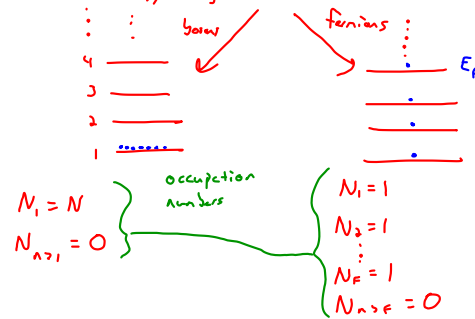
$$N \text{ (total states)} \propto k^3 \text{ in 3D} \Rightarrow N^{\frac{1}{3}} \propto k \Rightarrow E_{\text{tot}} \propto N^{\frac{2}{3}}$$

$$\propto L \text{ in 1D} \Rightarrow N \propto k \Rightarrow E_{\text{tot}} \propto N^2$$

difference in behavior between bosons and fermions

Quantum Statistical Mechanics

As $T \rightarrow 0$ a multiparticle system goes to lowest energy configuration



What happens when $T \neq 0$?

Basic postulate of statistical mechanics.

For a given E_{tot} , every ^{microscopic} configuration compatible w/ E_{tot} is equally likely.

So if the total # of micro. config. is γ then the probability of any one of those configurations is $P_i = \frac{1}{\gamma}$.

But if a set of configurations corresponds to the same set of occupation #'s, then $P_{set} = \frac{1}{\gamma} + \frac{1}{\gamma} + \dots$.

So the most probable set of occupation #'s corresponds to the set of N_n 's that can be realized in the most ways.

of microstates for a choice of occupation #'s N_n is $Q(N_1, N_2, \dots)$ $N = N_1 + N_2 + N_3 + N_4 + N_5$

\leftarrow Easy pp.

$$Q(0, 0, N, 0, 0) = 1$$

$$Q(N_1, N_2, N_3, N_4, N_5) \neq 1$$

Today: We will calculate Q 's

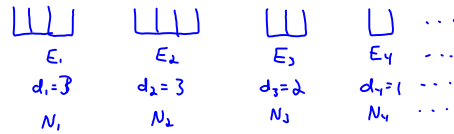
Next time: We extremize the Q 's to find most probable occupation #'s.

Counting (Combinatorics)

To calculate $Q(N_1, N_2, \dots)$ we start w/ a set of particles N .

Now choose a set of occupation numbers (N_1, N_2, \dots) comp. w/ $N = N_1 + N_2 + \dots$

The reason this is nontrivial is due to degeneracy.



- 2 step method:
1. Choose N_1 of the N total particles.
 2. Place N_1 particles in E_1 bins

Distinct Particles:

1. N ways to choose particle 1
- $N-1$ " " " " " 2
- $N-2$ " " " " " 3
- $N-N_i+1$ " " " " " N_i

$N(N-1)(N-2)\dots(N-N_i+1)$ ways to choose but they are ordered. To remove ordering divide by $N!$.

In total the number of ways to select N_i particles from N w/out ordering is $\frac{N(N-1)\dots(N-N_i+1)}{N!} = \frac{N!}{N!(N-N_i)!} \equiv \binom{N}{N_i}$

2. We have N_i particles to distribute to d_i bins.
 - Particle 1 can go into d_i bins.
 - Particle 2 " " " " d_i bins.
 - $d_i^{N_i}$ ways to distribute

Total # of ways to choose N_i particles and distribute among d_i bins is: $\binom{N}{N_i} d_i^{N_i}$

For E_2 , choose N_2 from $N-N_1$

$$\binom{N-N_1}{N_2} d_2^{N_2}$$

The total # of ways of populating N_1, N_2, \dots

$$Q(N_1, N_2, \dots) = \frac{N! d_1^{N_1}}{N_1!(N-N_1)!} \frac{(N-N_1)! d_2^{N_2}}{N_2!(N-N_1-N_2)!} \dots$$

$$= N! \prod_{n=1}^{\infty} \frac{d_n^{N_n}}{N_n!}$$

